

# Nonlinear Oscillations

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## I.) Nonlinear Oscillations - An Introduction

### a.) Overview / Roadmap

Can formulate sequence of problems:

- 1.) Simple harmonic oscillator - linear
- external force  $\left\{ \begin{array}{l} \text{deterministic} \\ \text{random} \end{array} \right.$

$$\ddot{x} + \alpha \dot{x} + \omega_0^2 x = f_{\text{ext}}$$

$\uparrow$  damping (friction const.)  $\rightarrow k/m$

if  $f_{\text{ext}} = f e^{i\gamma t} \Rightarrow \begin{cases} x = A e^{i\gamma t} \\ A = F/m(\omega_0^2 - \gamma^2 + 2i\alpha\gamma) \end{cases}$

- 2.) Variation: SHO (Linear), with Parametric Variation  
i.e.  $\omega_0^2 = \omega_0^2(t)$

$\omega_0^2(t)$   $\left\{ \begin{array}{l} \text{slow variation } \frac{1}{\omega_0^2} \frac{d}{dt} \omega_0^2 \ll \omega_0 \Rightarrow \text{adiabatic theory} \\ \text{fast variation } \frac{1}{\omega_0} \frac{d}{dt} \omega_0^2 > \omega_0 \Rightarrow \text{parametric instability} \\ \sim 2\omega_0 \end{array} \right.$

For parametric instability, prototype system is:

$$\begin{cases} \ddot{x} + \omega^2(t)x = 0 \\ \omega^2(t) = \omega_0^2(1+h \cos(\gamma t)) \\ \gamma \sim 2\omega_0 + \epsilon, \quad \epsilon \ll \omega_0 \end{cases} \begin{cases} \omega^2 = \omega^2(t) \\ \Rightarrow \partial L / \partial t \neq 0 \\ \Rightarrow \text{energy not const., here} \end{cases}$$

Mathieu's Eqn.  
 stability for  $\gamma \approx 2\omega_0$  resonance

then true nonlinear problems.

### 3.) Nonlinear Oscillator - Conservative

i.e.  $\ddot{x} + \omega_0^2 x + \epsilon x^3 = 0$  Duffing's Equation

↑  
anharmonic  
nonlinear } term

Here, conservative system, with bounded phase space trajectories

i.e.  $\frac{1}{2} \dot{x}^2 + F(x) = E$   $\rightarrow$  quasi-periodic orbits bounded

↓  
 $\omega_0^2 x^2 / 2 + \epsilon x^4 / 4$

$\rightarrow$  secularities in perturbation theory

$$\ddot{x} + \omega_0^2 x + \epsilon x^3 = 0$$

$$\omega^2 = \omega_0^2 \langle x^2 \rangle + \epsilon \langle x^4 \rangle \quad \downarrow \quad x^{(0)} \sim a \sin \omega t$$

$$\langle x^2 \rangle \sim a^2 \sin^2 \omega t \quad \epsilon x^3 \sim a^3 \sin^3 \omega t \sim a^3 \sin \omega t \quad (\text{on resonance})$$

$\Rightarrow$  secularities  $\sim t \cos \omega t$  ? ? ?

$\Rightarrow$  nonlinear frequency shift /  $\rightarrow$  removes secularities!

$\rightarrow$  simplest example of Reductive P.T.

$$\Rightarrow \omega = \omega_0 + \omega^{(2)}$$

amplitude dep.

i.e.  $\omega_0^2 x + \epsilon x^3$   
 $\rightarrow (\omega_0^2 + \epsilon \overline{x^2}) x$   
 $\sim (\omega_0^2 + (\#)\epsilon a^2) x$

$\uparrow$   
 2<sup>nd</sup>-order correction  $\sim \frac{3\epsilon a^2}{8\omega_0}$

$\hookrightarrow$  effective  $\omega$ -correction!

### 4.) Forced Nonlinear Oscillator

i.e.  $\ddot{x} + \alpha \dot{x} + \omega_0^2 x + \epsilon x^3 = f(t)$   
 $\uparrow$  NL  $\uparrow$  forcing

Forced Duffing's Equation

Here  $\rightarrow$  role of nonlinearity/anharmonicity in resonance

$\rightarrow$  i.e. contrast: forced SHO ( $\epsilon=0$ )

$f_{ext} = f e^{i\omega t}$

$x = B e^{i\omega t}$

dissipation resolves resonance  
 $\downarrow$

$B = b e^{i\phi}$

$b = f/m [(\omega_0^2 - \omega^2)^2 + 4\alpha^2 \omega^2]^{1/2}$

$\tan \phi = 2\alpha\omega / (\omega^2 - \omega_0^2)$

so:

$|b|^2 = \frac{f^2}{m^2} [(\omega_0^2 - \omega^2)^2 + 4\alpha^2 \omega^2] \rightarrow \frac{f^2}{4m^2 \omega_0^2} \left( \frac{1}{(\omega_0 - \omega)^2 + \alpha^2} \right)$   
 resonance  $\rightarrow \omega_0$   
 linewidth  $\rightarrow \alpha$   
 (symmetric)  
 (near resonance)

but for NL forced SHO;

→ nonlinear frequency shift enters resonance!  
here:

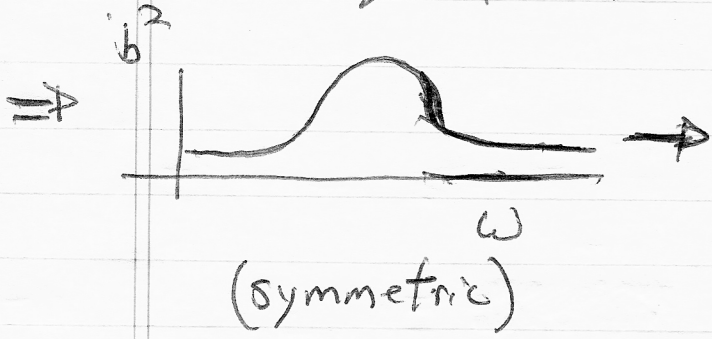
$$\omega - \omega_0 \rightarrow \omega - \omega_0 = Kb^2$$

{ amplitude-dependent  
shift

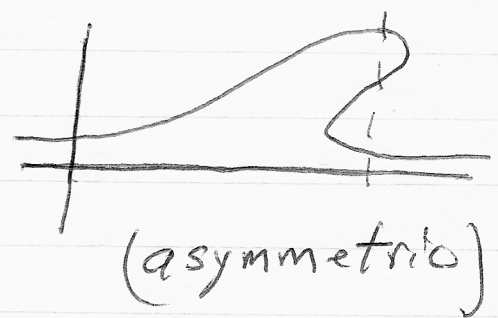
→

$$b^3 = \left( \frac{F}{4m^2\omega_0^2} \right) \frac{1}{[(\omega - \omega_0 - Kb^2)^2 + d^2]}$$

→ cubic eqn. for  $b^2$



(symmetric)



(asymmetric)

S-curve, bifurcations,  
mode jumping, etc.

### 6e) Nonlinear Oscillator with $\pm$ Dissipation

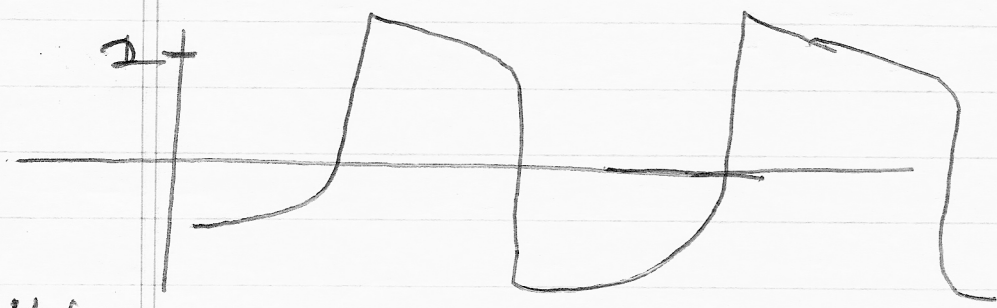
e.g. Van-der-Pol oscillator / Eqn.:

$$\ddot{x} + \alpha(x^2 - 1)\dot{x} + \omega_0^2 x = 0$$

key feature:  $\frac{\text{nonlinear}}{\text{variable sign}} > \frac{\text{dissipation}}{\text{(friction)}}$

$$\Rightarrow \begin{cases} x^2 < 1 \rightarrow \alpha(x^2 - 1) < 0 \Rightarrow \begin{cases} \text{negative dissipation} \\ \text{instability} \end{cases} \\ x^2 > 1 \rightarrow \alpha(x^2 - 1) > 0 \Rightarrow \begin{cases} \text{positive dissipation} \\ \rightarrow \text{stability / saturation} \end{cases} \end{cases}$$

$\Rightarrow$  bursty behavior, characteristic of limit cycle (at  $x^2 = 1$ )



Note:

In contrast to previous examples, system is fundamentally dissipative - nonlinearity in friction.

System/Eqn.	Nonlinearity	External
Mathieu (Parametric)	None	$\omega_0^2 = \omega_0^2(t)$ Fast
Duffing (Freq. Shift)	Potential	None
Forced Duffing (Bi-stability/Bifurcation)	Potential	Forcing
Van-der-Pol (Limit Cycle)	Friction	None